

# Curved spacetime and implications for cosmology

- General Relativity implies spacetime is curved in the presence of matter
  - ▶ since universe contains matter, might expect overall curvature (as well as local “gravity wells”)
  - ▶ how does this affect measurements of large-scale distances?
  - ▶ what are the implications for cosmology?

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## Curved spacetime

- Two-dimensional curved space: surface of sphere

▶ distance between  $(r, \theta)$  and  $(r+dr, \theta+d\theta)$  given by

$$\text{▶ } ds^2 = dr^2 + R^2 \sin^2(r/R) d\theta^2$$

▶  $r$  = distance from pole  
 $\theta$  = angle from meridian  
 $R$  = radius of sphere

▶ positive curvature



Universe with positive curvature. Lines converge at great distances. Triangle angles add to more than 180°.

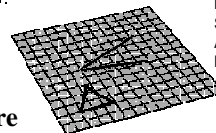


Universe with negative curvature. Lines diverge at increasing angles. Triangle angles add to less than 180°.

- “Saddle” (negative curvature)

▶  $ds^2 = dr^2 + R^2 \sinh^2(r/R) d\theta^2$

▶ (2D surface of constant negative curvature can't really be constructed in 3D space)



Universe with no curvature. Lines diverge at constant angle. Triangle angles add to 180°.

Nick Strobel's Astronomy Notes

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## ***3D curved spacetime***

- **Robertson-Walker metric**

- ▶  $ds^2 = -c^2 dt^2 + a^2(t)[dx^2/(1 - kx^2/R^2) + x^2(d\theta^2 + \sin^2\theta d\phi^2)]$ 
  - ▶ note sign change from our previous definition of  $ds^2$ !
- ▶  $a(t)$  is an overall scale factor allowing for expansion or contraction ( $a(t_0) \equiv 1$ )
- ▶  $x$  is called a comoving coordinate (unchanged by overall expansion or contraction)
- ▶  $k$  defines sign of curvature ( $k = \pm 1$  or  $0$ ),  $R$  is radius of curvature
- ▶ path of photon has  $ds^2 = 0$ , as before

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## ***Implications for cosmology***

- comoving proper distance ( $dt = 0$ ) between origin and object at coordinate  $x$ :

$$r = \int_0^x \frac{dx}{\sqrt{1 - kx^2/R^2}}$$

- ▶ for  $k = +1$  this gives  $r = R \sin^{-1}(x/R)$ , i.e.  $r \leq 2\pi R$ 
  - ▶ finite but unbounded universe, cf. sphere
- ▶ for  $k = -1$  we get  $r = R \sinh^{-1}(x/R)$ , and for  $k = 0$ ,  $r = x$ 
  - ▶ infinite universe, cf. saddle
- ▶ for  $x \ll R$  all values of  $k$  give  $r \approx x$ 
  - ▶ any spacetime looks flat on small enough scales
- ▶ this is independent of  $a$ 
  - ▶ it's a comoving distance

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## Implications for cosmology

- cosmological redshift: change variables in RW metric from  $x$  to  $r$ :

- ▶  $ds^2 = -c^2 dt^2 + a^2(t)[dr^2 + x(r)^2 d\Omega^2]$

- ▶ for light  $ds = 0$ , so  $c^2 dt^2 = a^2(t) dr^2$ , i.e.  $c dt/a(t) = dr$  (assuming beam directed radially)

- ▶ suppose wave crest emitted at time  $t_e$  and observed at  $t_o$

$$c \int_{t_e}^{t_o} \frac{dt}{a(t)} = \int_0^r dr = r$$

first wave crest

$$c \int_{t_e + \lambda_e/c}^{t_o + \lambda_o/c} \frac{dt}{a(t)} = \int_0^r dr = r$$

next wave crest

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## Implications for cosmology

- Then 
$$c \int_{t_e}^{t_e + \lambda_e/c} \frac{dt}{a(t)} = c \int_{t_o}^{t_o + \lambda_o/c} \frac{dt}{a(t)}$$

but if  $\lambda \ll c/H_0$ ,  $a(t)$  is almost constant over this integral, so we can write

$$\frac{c}{a(t_e)} \int_{t_e}^{t_e + \lambda_e/c} dt = \frac{c}{a(t_o)} \int_{t_o}^{t_o + \lambda_o/c} dt$$

i.e. 
$$\frac{\lambda_e}{a(t_e)} = \frac{\lambda_o}{a(t_o)}$$

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## ***Implications for cosmology***

- So expanding universe produces redshift  $z$ , where

$$1 + z = \frac{a(t_o)}{a(t_e)} = \frac{1}{a(t_e)}$$

- Note:

- ▶  $z$  can have any value from 0 to  $\infty$
- ▶  $z$  is a measure of  $t_e$
- ▶ often interpret  $z$  using relativistic Doppler shift formula

$$1 + z = \sqrt{\frac{c + v}{c - v}}$$

but note that this is misleading: the object is **not** changing its local coordinates

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## ***Implications for cosmology***

- Conclusions

- ▶ in general relativity universe can be infinite (if  $k = -1$  or 0) or finite but unbounded (if  $k = +1$ )
- ▶ universe can expand or contract (if overall scale factor  $a(t)$  is not constant)
- ▶ if universe expands or contracts, radiation emitted by a comoving source will appear redshifted or blueshifted respectively

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