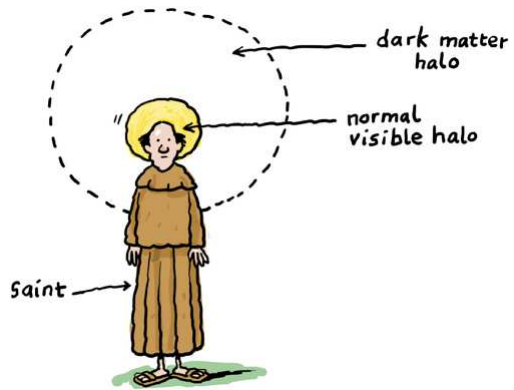




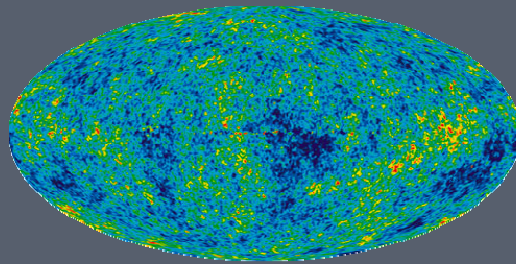
The University Of Sheffield.



1

NON-BARYONIC DARK MATTER

Susan Cartwright
University of Sheffield



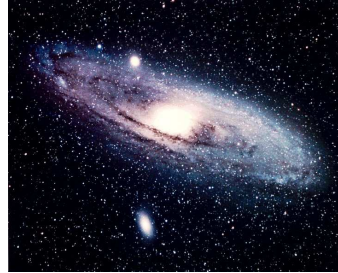
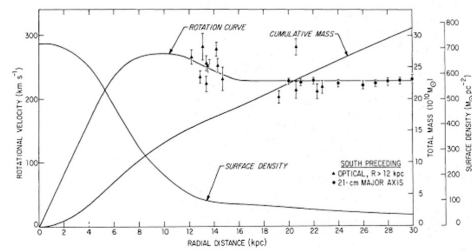
2

DARK MATTER

Astrophysical Evidence
Candidates
Detection

THE ASTROPHYSICAL EVIDENCE

○ Rotation curves of spiral galaxies



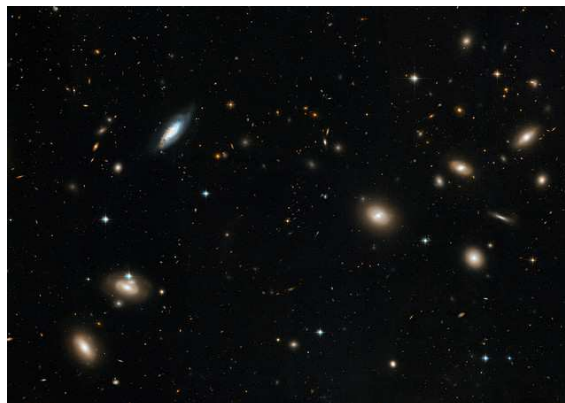
- flat at large radii: if mass traced light we would expect them to be Keplerian at large radii, $v \propto r^{-1/2}$, because the light is concentrated in the central bulge
 - and disc light falls off exponentially, not $\propto r^{-2}$ as required for flat rotation curve

3

THE ASTROPHYSICAL EVIDENCE

○ Dynamics of rich clusters

- Zwicky (1933!) noted that the velocities of galaxies in the Coma cluster were too high to be consistent with a bound system

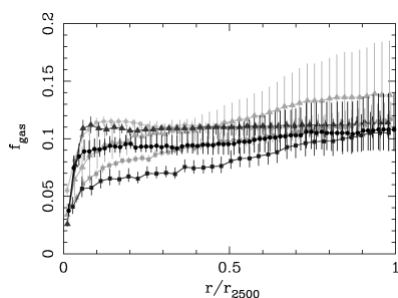


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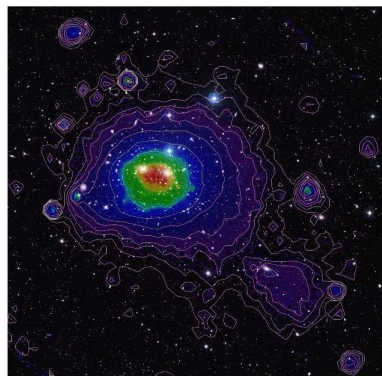
THE ASTROPHYSICAL EVIDENCE

○ Dynamics of rich clusters

- mass of gas and gravitating mass can be extracted from X-ray emission from intracluster medium



Allen et al., *MNRAS* **334** (2002) L11

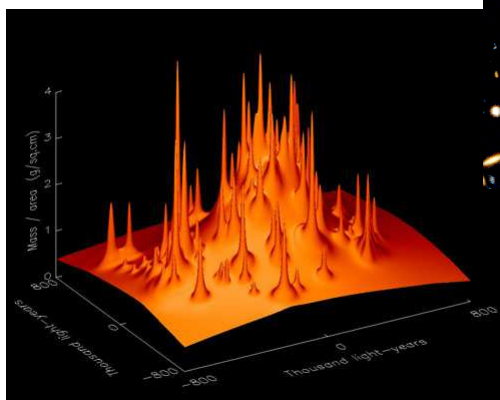


ROSAT X-ray image of Coma cluster overlaid on optical.
MPI (ROSAT image);
NASA/ESA/DSS2 (visible image)

5

THE ASTROPHYSICAL EVIDENCE

○ Dynamics of rich clusters



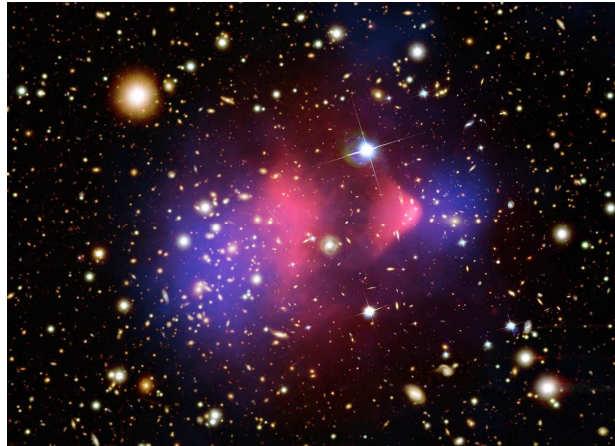
Mass map of CL0024+1654 as determined from the observed gravitational lensing.
Tyson, Kochanski and Dell'Antonio,
ApJ **498** (1998) L107

6

THE ASTROPHYSICAL EVIDENCE: THE BULLET CLUSTER

- Mass from lens mapping (blue) follows stars not gas (red)
 - dark matter is **collisionless**

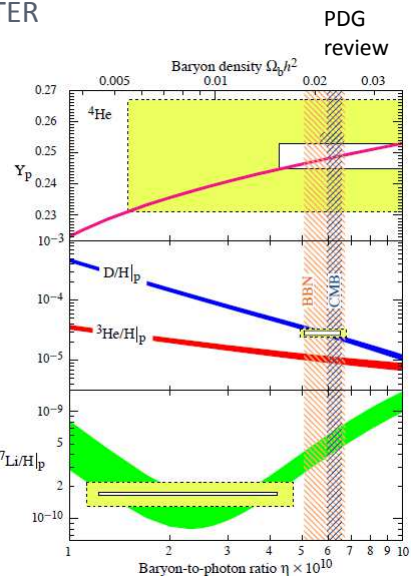
Composite Credit:
 X-ray: NASA/CXC/CfA/
 M. Markevitch et al.;
 Lensing Map:
 NASA/STScI; ESO WFI;
 Magellan/U.Arizona/
 D.Clowe et al
 Optical: NASA/STScI;
 Magellan/U.Arizona/
 D.Clowe et al.



7

NON-BARYONIC DARK MATTER

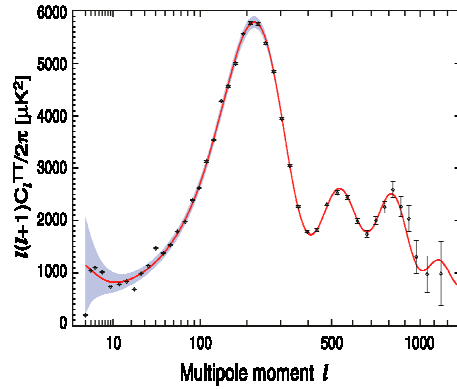
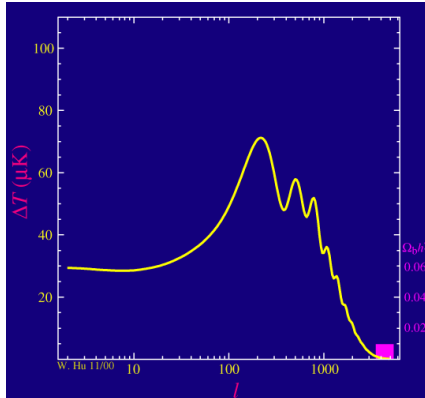
- Density of baryonic matter strongly constrained by early-universe nucleosynthesis (BBN)
 - density parameter of order 0.3 as required by data from, e.g., galaxy clusters is completely inconsistent with best fit
 - note here cosmologists' $h = H_0/100 \text{ km s}^{-1} \text{ Mpc}^{-1} \approx 0.7$
 - do not confuse with Planck's constant!



8

NON-BARYONIC DARK MATTER: COSMOLOGY

Ratio of odd/even peaks depends on Ω_b



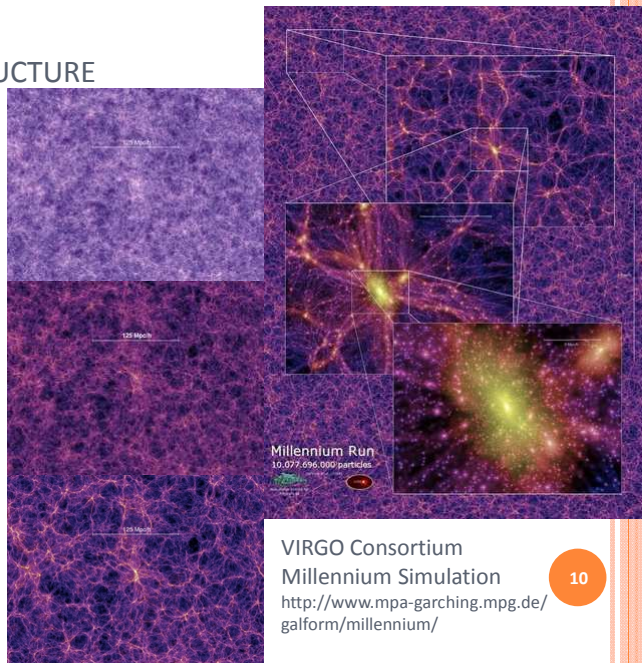
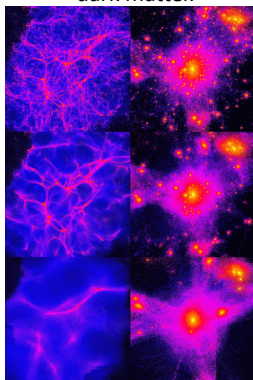
Wayne Hu

9

LARGE SCALE STRUCTURE

Relativistic (**hot**) dark matter makes structure top-down—non-relativistic (**cold**) bottom-up.

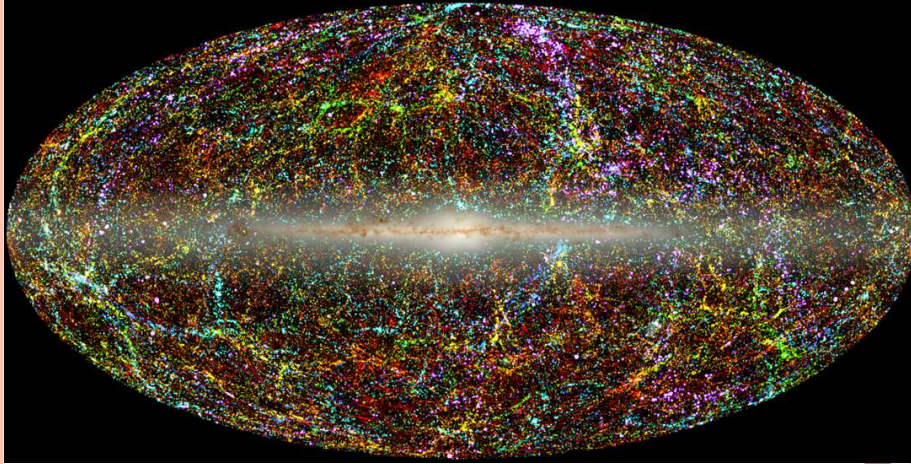
Real world looks like **cold** dark matter.



VIRGO Consortium
Millennium Simulation
<http://www.mpa-garching.mpg.de/galform/millennium/>

10

2MASS GALAXY SURVEY



Local galaxies ($z < 0.1$; distance coded by colour, from blue to red)

Statistical studies, e.g. correlation functions, confirm visual impression that this looks much more like cold than hot dark matter

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BRIEF SUMMARY OF ASTROPHYSICAL EVIDENCE

- Many observables concur that $\Omega_{m0} \approx 0.3$

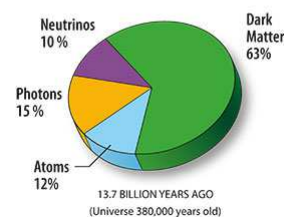
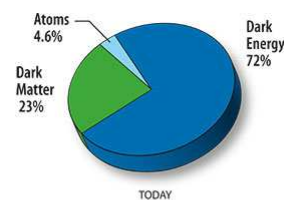
- Most of this must be non-baryonic

- BBN and CMB concur that baryonic matter contributes $\Omega_{b0} \approx 0.05$
- Bullet Cluster mass distribution indicates that dark matter is collisionless

- No Standard Model candidate

- neutrinos are too light, and are "hot" (relativistic at decoupling)
 - hot dark matter does not reproduce observed large-scale structure

→ *BSM physics*



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DARK MATTER

Astrophysical Evidence

Candidates

Detection

DARK MATTER CANDIDATES

	WIMPs	SuperWIMPs	Light \tilde{G}	Hidden DM	Sterile ν	Axions
Motivation	GHP	GHP	GHP/NPFP	GHP/NPFP	ν Mass	Strong CP
Naturally Correct Ω	Yes	Yes	No	Possible	No	No
Production Mechanism	Freeze Out	Decay	Thermal	Various	Various	Various
Mass Range	GeV-TeV	GeV-TeV	eV-keV	GeV-TeV	keV	μeV -meV
Temperature	Cold	Cold/Warm	Cold/Warm	Cold/Warm	Warm	Cold
Collisional				✓		
Early Universe		✓✓		✓		
Direct Detection	✓✓			✓		✓✓
Indirect Detection	✓✓	✓		✓	✓✓	
Particle Colliders	✓✓	✓✓	✓✓	✓		

GHP = Gauge Hierarchy Problem; NPFP = New Physics Flavour Problem
 ν = possible signal; $\nu\nu$ = expected signal

Jonathan Feng, *ARAA* **48** (2010) 495 (highly recommended)

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PARTICLE PHYSICS MOTIVATIONS

o Gauge Hierarchy Problem

- in SM, loop corrections to Higgs mass give

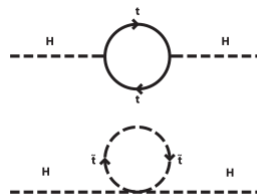
$$\Delta m_h^2 \approx \frac{\lambda^2}{16\pi^2} \int \frac{d^4 p}{p^2} \approx \frac{\lambda^2}{16\pi^2} \Lambda^2$$

and there is no obvious reason why $\Lambda \neq M_{Pl}$

- (Planck mass $M_{Pl} = (\hbar c/G)^{1/2} \approx 1.2 \times 10^{19}$ GeV = mass scale for quantum gravity)
- supersymmetry fixes this by introducing a new set of loop corrections that cancel those from the SM
- new physics at TeV scale will also fix it (can set $\Lambda \sim 1$ TeV)

o New Physics Flavour Problem

- we observe conservation or near-conservation of B, L, CP
 - and do not observe flavour-changing neutral currents
- new physics has a nasty tendency to violate these
 - can require fine-tuning or new discrete symmetries, e.g. *R*-parity

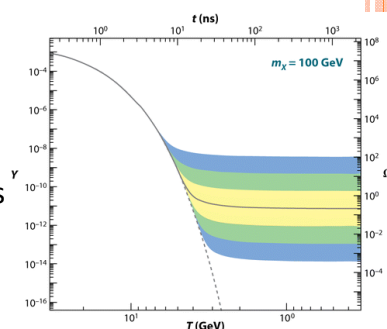


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WIMPs

o Weakly Interacting Massive Particles

- produced thermally in early universe
- annihilate as universe cools, but “freeze out” when density drops so low that annihilation no longer occurs with meaningful rate
 - “target volume” per particle in time Δt is $\sigma_A v \Delta t$, where σ_A is cross-section
 - so annihilation rate is $n_f \langle \sigma_A v \rangle$ where n_f is number density
- freeze-out occurs when $H \approx n_f \langle \sigma_A v \rangle$, and in radiation era we have $H \propto T^2/M_{Pl}$ (because $\rho \propto T^4$ and $G \propto 1/M_{Pl}^2$)
- can estimate relic density by considering freeze-out



Feng JL. 2010. Annu. Rev. Astron. Astrophys. 48:495–545

$$n_f \approx (m_\chi T_f)^{3/2} e^{-m_\chi/T_f} \approx \frac{T_f^2}{M_{Pl} \langle \sigma_A v \rangle}$$

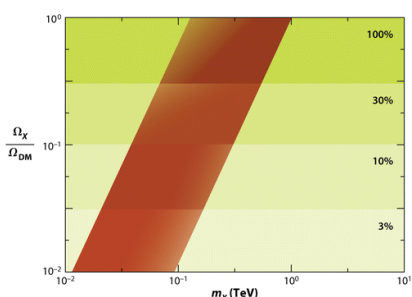
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WIMP RELIC DENSITY

○ Converting to Ω gives $\Omega_x = \frac{m_x n_0}{\rho_c} \approx \frac{m_x T_0^3 n_f}{\rho_c T_f^3} \approx \frac{x_f T_0^3}{\rho_c M_{Pl}} \langle \sigma_{AV} \rangle^{-1}$
 where $x_f = m_x/T_f$

- and typically $\langle \sigma_{AV} \rangle \propto 1/m_x^2$ or v^2/m_x^2 (S or P wave respectively)

○ Consequence: weakly interacting massive particles with electroweak-scale masses “naturally” have reasonable relic densities



- and therefore make excellent dark matter candidates

Feng JL. 2010. Annu. Rev. Astron. Astrophys. 48:495–545

SUPERSYMMETRY (SUSY)

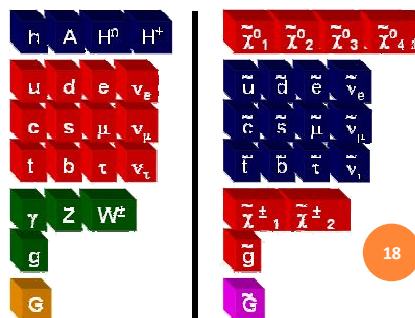
- Extension to Standard Model in which all fermions have partner particles that are bosons, and vice versa
 - if this were an exact symmetry we’d see twice as many particles
 - therefore it is a “broken” symmetry—sparticles much more massive than SM particles

○ Slightly extended “normal” particle content

- need to generate SUSY masses leads to extra Higgs particles

○ Some SUSY particles are mixed states

- neutralinos χ are mixed partners of Z, γ , h and H



SUPERSYMMETRIC WIMPS

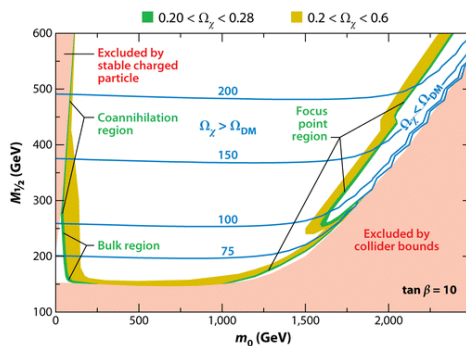
- Supersymmetry solves the GHP by introducing cancelling corrections
 - predicts a complete set of new particles
 - well-defined interactions, but unknown masses (10 GeV – few TeV)
 - NFP often solved by introducing *R*-parity—new discrete quantum number
 - then lightest supersymmetric particle is stable
 - best DM candidate is lightest neutralino (mixed partner of W^0 , B, H, h)
 - far too many free parameters in most general supersymmetric models
 - so usually consider constrained models with simplifying assumptions
 - most common constrained model: mSUGRA
 - parameters $m_0, M_{1/2}, A_0, \tan \beta, \text{sign}(\mu)$
 - mSUGRA neutralino is probably the best studied DM candidate

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SUSY WIMPS

this means that
particle = antiparticle

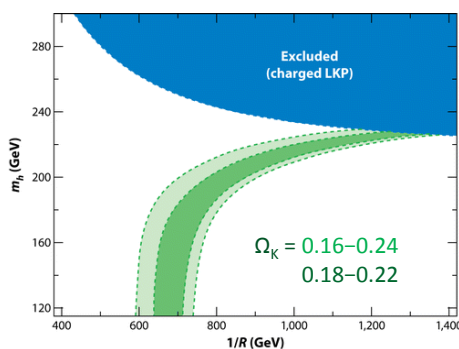
- Neutralinos are Majorana fermions and therefore self-annihilate
 - Pauli exclusion principle implies that $\chi_1\chi_1$ annihilation prefers to go to spin 0 final state
 - $f\bar{f}$ prefers spin 1
 - therefore annihilation cross-section is suppressed
 - hence Ω_χ tends to be too high
 - parameter space **very** constrained by WMAP



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KALUZA-KLEIN WIMPS

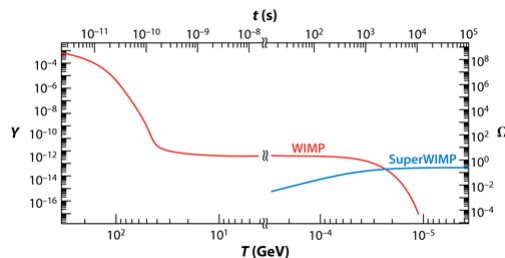
- In extra-dimension models, SM particles have partners with the *same* spin
 - “tower” of masses separated by R^{-1} , where R is size of compactified extra dimension
 - new discrete quantum number, K -parity, implies lightest KK particle is stable
 - this is the potential WIMP candidate
 - usually B^1
 - annihilation not spin-suppressed (it’s a boson), so preferred mass higher



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SUPERWIMPS

- Massive particles with superweak interactions
 - produced by decay of metastable WIMP
 - because this decay is superweak, lifetime is very long (10^3 – 10^7 s)
 - WIMP may be neutralino, but could be charged particle
 - dramatic signature at LHC (stable supermassive particle)
 - candidates:
 - weak-scale gravitino
 - axino
 - equivalent states in KK theories
- these particles cannot be directly detected, but indirect-detection searches and colliders may see them
 - they may also have detectable astrophysical signatures



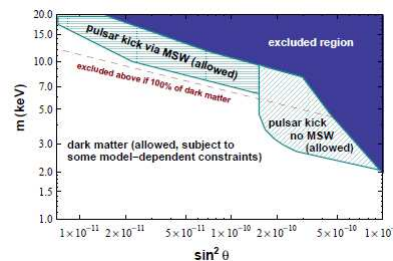
22

LIGHT GRAVITINOS

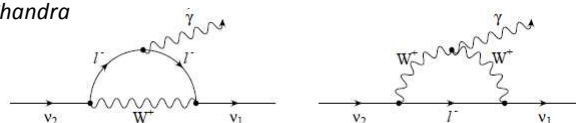
- Expected in gauge-mediated supersymmetry breaking
 - in these models gravitino has $m < 1$ GeV
 - neutralinos decay through $\tilde{\nu}\tilde{G}$, so cannot be dark matter
 - gravitinos themselves are possible DM candidates
 - but tend to be too light, i.e. too warm, or too abundant
 - relic density in minimal scenario is $\Omega_{\tilde{G}} \approx 0.25 m_{\tilde{G}}/(100 \text{ eV})$
 - so require $m_{\tilde{G}} < 100 \text{ eV}$ for appropriate relic density
 - but require $m_{\tilde{G}} > 2 \text{ keV}$ for appropriate large-scale structure
 - models which avoid these problems look contrived

STERILE NEUTRINOS

Kusenko, DM10



- Seesaw mechanism for generating small ν_L masses implies existence of massive right-handed sterile states
 - usually assumed that $M_R \approx M_{GUT}$, in which case sterile neutrinos are not viable dark matter candidates
 - but smaller Yukawa couplings can combine with smaller M_R to produce observed ν_L properties together with sterile neutrino at keV mass scale—viable dark matter candidate
 - such a sterile neutrino could also explain observed high velocities of pulsars (asymmetry in supernova explosion generating “kick”)
 - these neutrinos are not entirely stable: $\tau \gg 1/H_0$, but they do decay and can generate X-rays via loop diagrams—therefore potentially detectable by, e.g., *Chandra*



STERILE NEUTRINOS

○ Production mechanisms

- oscillation at $T \approx 100$ MeV
 - $\Omega_\nu \propto \sin^2 2\theta m^{1.8}$ from numerical studies
 - always present: requires small mass and very small mixing angle
 - not theoretically motivated: some fine tuning therefore required
- resonant neutrino oscillations
 - if universe has significant lepton number asymmetry, $L > 0$
- decays of heavy particles
 - e.g. singlet Higgs driving sterile neutrino mass term

○ Observational constraints

- X-ray background
- presence of small-scale structure
 - sterile neutrinos are “warm dark matter” with Mpc free-streaming

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AXIONS

○ Introduced to solve the “strong CP problem”

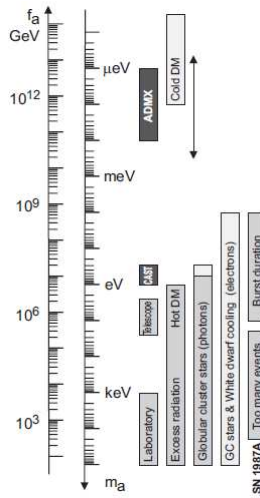
- SM Lagrangian includes CP-violating term which should contribute to, e.g., neutron electric dipole moment
 - neutron doesn't appear to *have* an EDM ($< 3 \times 10^{-26}$ e cm, cf. naïve expectation of 10^{-16}) so this term is strongly suppressed
- introduce new pseudoscalar field to kill this term (Peccei-Quinn mechanism)
 - result is an associated pseudoscalar boson, the axion

○ Axions are extremely light (< 10 meV), but are cold dark matter

- not produced thermally, but via phase transition in very early universe
 - if this occurs before inflation, visible universe is all in single domain
 - if after inflation, there are many domains, and topological defects such as axion domain walls and axionic strings may occur

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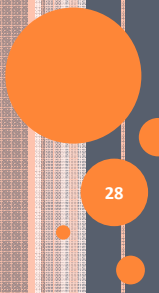
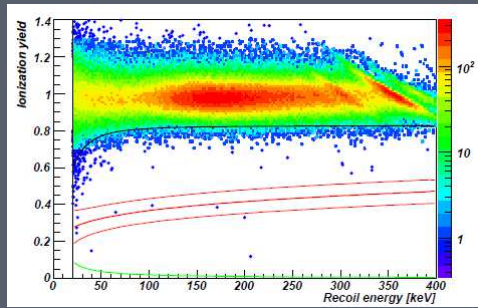
AXIONS



Georg Raffelt, hep-ph/0611350v1

- Axion mass is $m_a \approx 6 \mu\text{eV} \times f_a / (10^{12} \text{ GeV})$ where f_a is the unknown mass scale of the PQ mechanism
- Calculated relic density is $\Omega_a \approx 0.4 \vartheta^2 (f_a / 10^{12} \text{ GeV})^{1.18}$ where ϑ is initial vacuum misalignment
 - so need $f_a < 10^{12} \text{ GeV}$ to avoid overclosing universe
 - astrophysical constraints require $f_a > 10^9 \text{ GeV}$
 - therefore $6 \mu\text{eV} < m_a < 6 \text{ meV}$

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DARK MATTER

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- Astrophysical Evidence
- Candidates
- Detection

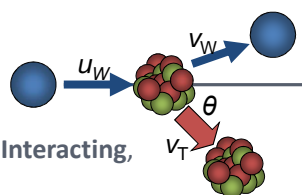
DETECTION OF DARK MATTER CANDIDATES

- Direct detection
 - dark matter particle interacts in your detector and you observe it
- Indirect detection
 - you detect its decay/annihilation products or other associated phenomena
- Collider phenomenology
 - it can be produced at, say, LHC and has a detectable signature
- Cosmology
 - it has a noticeable and characteristic impact on BBN or CMB
- Focus here on best studied candidates—WIMPs and axions

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DIRECT DETECTION: WIMP-NUCLEUS INTERACTION

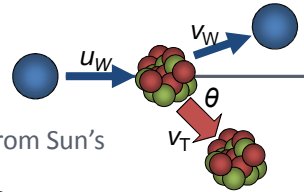
- Key points:
 - it doesn't happen very often: **Weakly Interacting**, remember?
 - it is **non-relativistic**: WIMPs are bound in Galactic halo, so have velocities ~ 220 km/s ($v/c \sim 10^{-3}$)
 - it is **elastic scattering**—momentum and KE conserved
- If we assume that spin plays no role, we can model this as collision of two hard spheres of masses M_W, M_T
 - we find that $v_T = \frac{2M_W}{M_W + M_T} u_W \cos \theta$
 - assuming nucleus initially at rest, $u_T = 0$
 - maximal for head-on scattering ($\cos \theta = 1$), and for $M_W = M_T$
 - u_W and its likely direction can be calculated by modelling the halo



de Broglie wavelength of particle with 1 TeV mass is $h/p \approx 10^{-15}$ m \approx nuclear radius

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WIMP-NUCLEUS INTERACTION



Basic numbers:

- local density of DM can be deduced from Sun's orbital velocity via

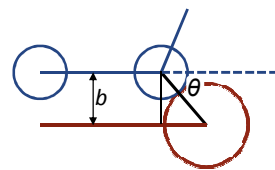
$$\rho(R_{Sun}) = \frac{1}{4\pi R_{Sun}^2} \frac{dM_r}{dr} \Big|_{R_{Sun}} = \frac{1}{4\pi R_{Sun}^2} \frac{V^2}{G}$$
- this gives 0.3–0.5 GeV/cm³ depending on exactly what you assume for V and R_{sun} (neither of which is very well known)
- WIMP rest energy expected to be in range 10–1000 GeV
 - so, between 0.3 and 50 particles per litre in solar neighbourhood
 - note that this assumes halo is an isothermal sphere—it might not be!
- Kinetic energy of WIMP $\frac{1}{2}M_W V^2 \approx 2.7\text{--}270$ keV if $V \sim 220$ km/s
 - best case scenario: **all** of this transferred to nucleus—but this will not normally happen (requires $\cos \theta = 1$ and $M_W = M_T$)

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WIMP-NUCLEUS INTERACTION: ENERGY SPECTRUM

Scattering angle depends on impact parameter b

- $\sin \theta = b / (R_W + R_T) = b / R$
- Probability of impact parameter between b and $b + db$ is area of shaded region divided by total area = $2\pi b db / \pi R^2 = (2b/R^2) db$

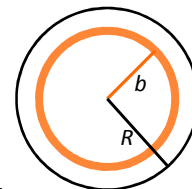


Transferred energy is $\frac{1}{2}M_T v_T^2 = E_T$ where

$$E_T = \frac{M_T M_W}{(M_T + M_W)^2} E_W \cos^2 \theta$$

- $P(\cos^2 \theta) = P(b) / |d(\cos^2 \theta)/db| = 1$
- All values of recoil energy are equally likely

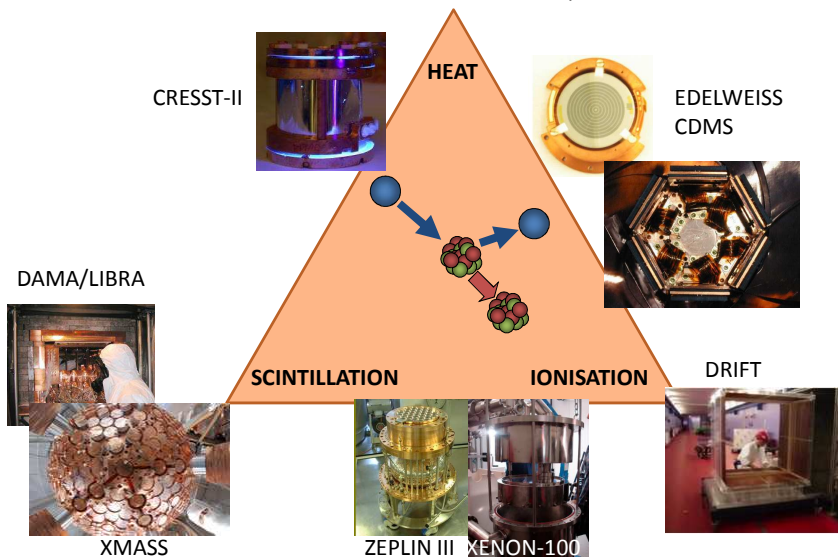
- and for a given halo model the only unknown is M_W



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DIRECT DETECTION OF WIMPS

Basic principle: WIMP scatters elastically from nucleus; experiment detects nuclear recoil



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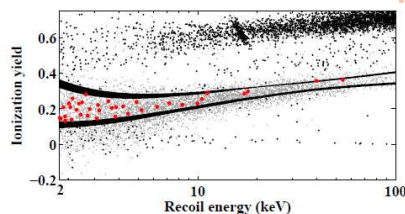
DIRECT DETECTION OF WIMPS

Backgrounds

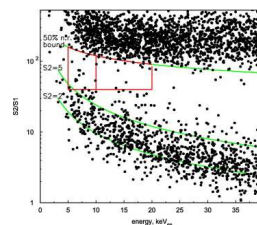
- cosmics and radioactive nuclei (especially radon)
 - use deep site and radiopure materials
 - use discriminators to separate signal and background

Time variation

- expect annual variation caused by Earth's and Sun's orbital motion
 - small effect, ~7%
 - basis of claimed signal by DAMA experiment
- much stronger diurnal variation caused by changing orientation of Earth
 - "smoking gun", but requires directional detector
 - current directional detector, DRIFT, has rather small target mass (being gaseous)—hence not at leading edge of sensitivity



CDMS-II, PRL 106 (2011) 131302

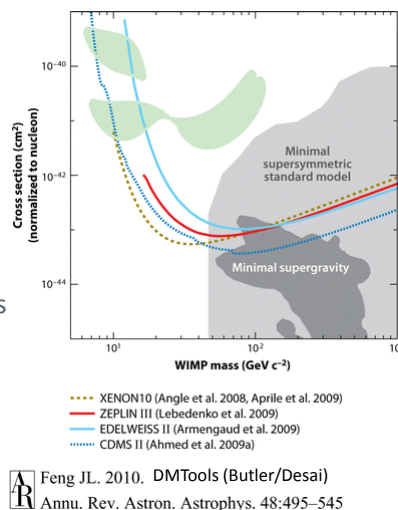


ZEPLIN-II, Astropart. Phys. 28 (2007) 287

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DIRECT DETECTION OF WIMPS

- Interaction with nuclei can be spin-independent or spin-dependent
 - spin-dependent interactions require nucleus with net spin
 - most direct detection experiments focus on SI, and limits are much better in this case
- Conflict between DAMA and others tricky to resolve
 - requires very low mass and high cross-section
 - if real, may point to a non-supersymmetric DM candidate



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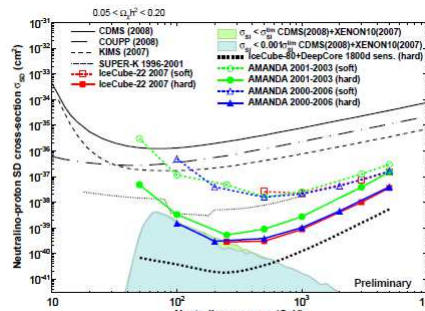
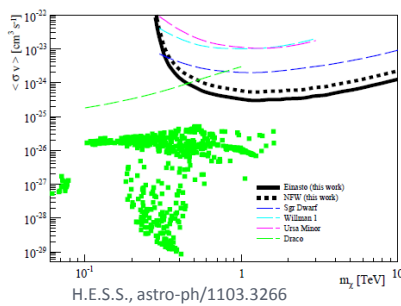
INDIRECT DETECTION OF WIMPS

- After freeze-out, neutralino self-annihilation is negligible in universe at large
 - but neutralinos can be captured by repeated scattering in massive bodies, e.g. Sun, and this will produce a significant annihilation rate
 - number of captured neutralinos $N = C - AN^2$ where C is capture rate and A is $\langle \sigma_A v \rangle$ per volume
 - if steady state reached, annihilation rate is just $C/2$, therefore determined by scattering cross-section
 - annihilation channels include W^+W^- , $b\bar{b}$, $\tau^+\tau^-$, etc. which produce secondary neutrinos
 - these escape the massive object and are detectable by neutrino telescopes

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INDIRECT DETECTION OF WIMPS

- Relatively high threshold of neutrino telescopes implies greater sensitivity to “hard” neutrinos, e.g. from WW
- Also possible that neutralinos might collect near Galactic centre
 - in this region other annihilation products, e.g. γ -rays, could escape



- search by H.E.S.S. found nothing
- signals at lower energies could be astrophysical not astroparticle

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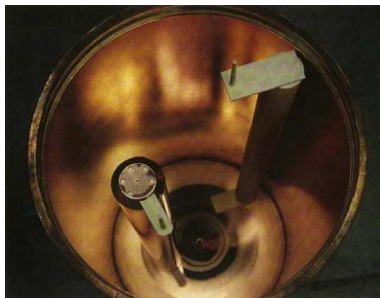
LHC DETECTION OF WIMPS AND SWIMPS

- WIMPs show up at LHC through missing-energy signature
 - note: not immediate proof of dark-matter status
 - long-lived but not stable neutral particle would have this signature but would not be DM candidate
 - need to constrain properties enough to calculate expected relic density if particle *is* stable, then check consistency
- SuperWIMP parents could also be detected
 - if charged these would be spectacular, because of extremely long lifetime
 - very heavy particle exits detector without decaying
 - if seen, could in principle be trapped in external water tanks, or even dug out of cavern walls (Feng: “new meaning to the phrase ‘data mining’”)
 - if neutral, hard to tell from WIMP proper
 - but mismatch in relic density, or conflict with direct detection, possible clues

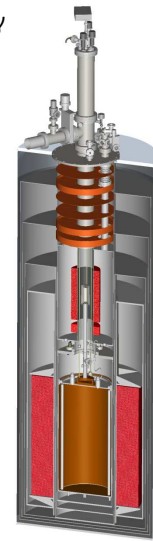
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AXION DETECTION

- Axions couple (unenthusiastically) to photons via $\mathcal{L}_{a\gamma\gamma} = -g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$
 - they can therefore be detected using Primakoff effect (resonant conversion of axion to photon in magnetic field)
 - ADMX experiment uses very high Q resonant cavity in superconducting magnet to look for excess power

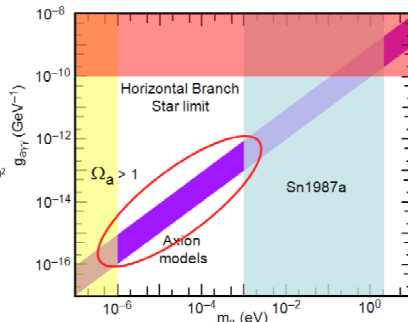
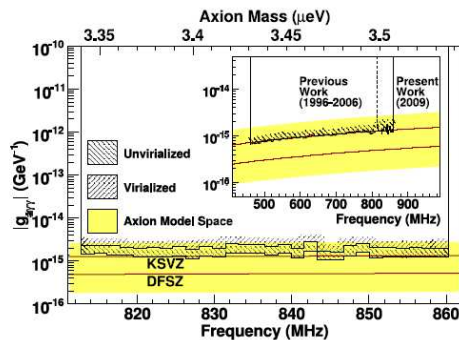
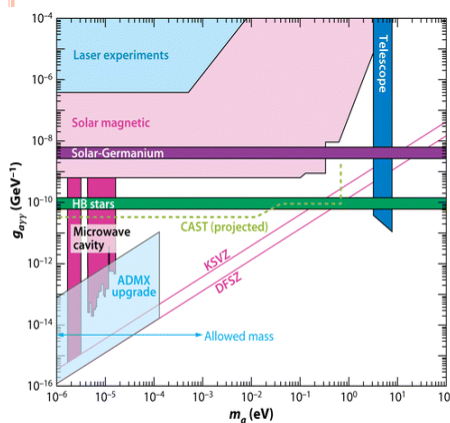


- this is a scanning experiment: need to adjust resonant frequency to "see" specific mass (very tedious)
- alternative: look for axions produced in Sun (CAST)
 - non-scanning, but less sensitive




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AXION DETECTION



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 Feng JL. 2010. Annu. Rev. Astron. Astrophys. 48:495–545

DARK MATTER: SUMMARY

- Astrophysical evidence for dark matter is consistent and compelling
 - not an unfalsifiable theory—for example, severe conflict between BBN and WMAP on Ω_b might have scuppered it
- Particle physics candidates are many and varied
 - and in many cases are not *ad hoc* inventions, but have strong independent motivation from within particle physics
- Unambiguous detection is possible for several candidates, but will need careful confirmation
 - interdisciplinary approaches combining direct detection, indirect detection, conventional high-energy physics and astrophysics may well be required

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THE END



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