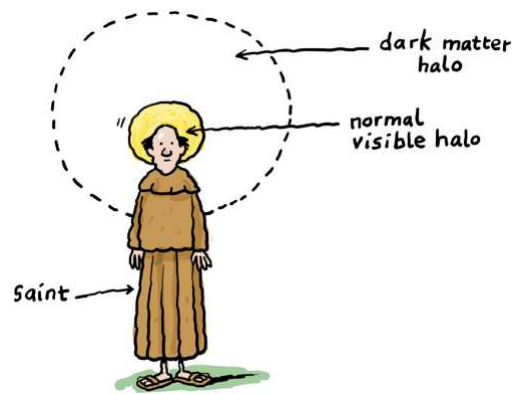
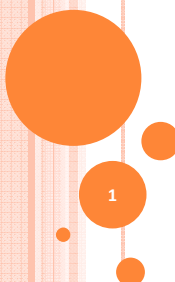


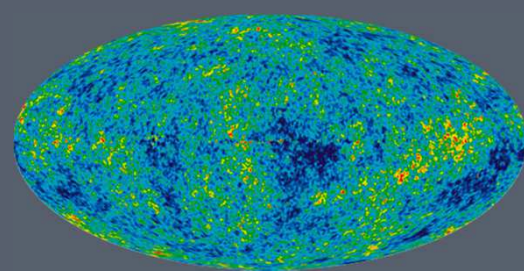
The University  
Of  
Sheffield.

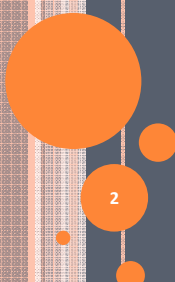




## ASTROPARTICLE PHYSICS LECTURE 4

1  
Susan Cartwright  
University of Sheffield



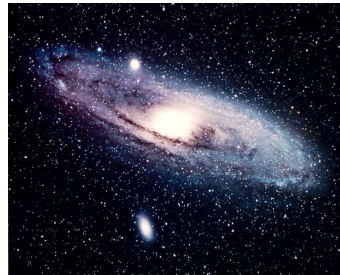
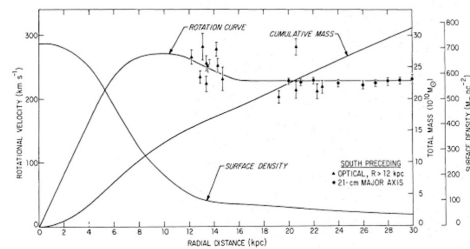


## DARK MATTER

2  
Astrophysical Evidence  
Candidates  
Detection

## THE ASTROPHYSICAL EVIDENCE

### Rotation curves of spiral galaxies



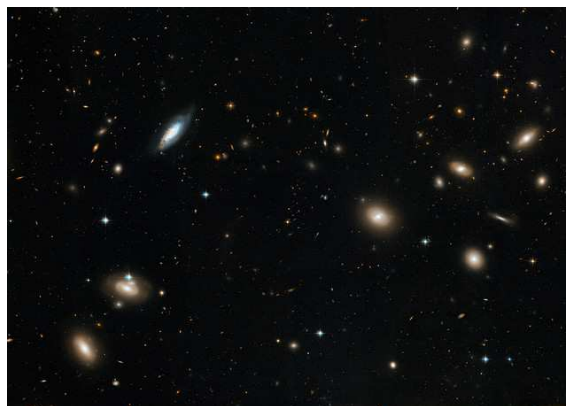
- flat at large radii: if mass traced light we would expect them to be Keplerian at large radii,  $v \propto r^{-1/2}$ , because the light is concentrated in the central bulge
  - and disc light falls off exponentially, not  $\propto r^{-2}$  as required for flat rotation curve

3

## THE ASTROPHYSICAL EVIDENCE

### Dynamics of rich clusters

- Zwicky (1933!) noted that the velocities of galaxies in the Coma cluster were too high to be consistent with a bound system

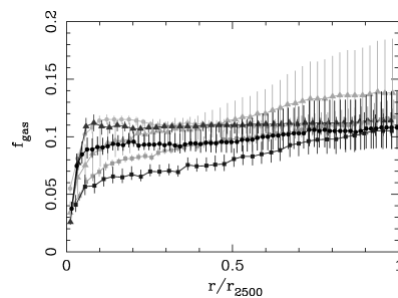


4

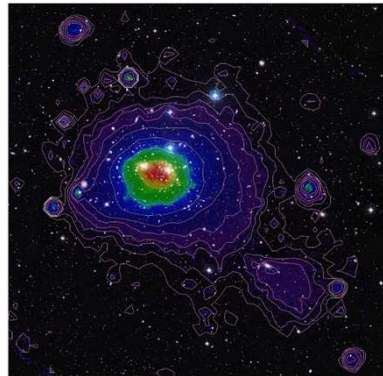
## THE ASTROPHYSICAL EVIDENCE

### ○ Dynamics of rich clusters

- mass of gas and gravitating mass can be extracted from X-ray emission from intracluster medium



Allen et al., *MNRAS* **334** (2002) L11

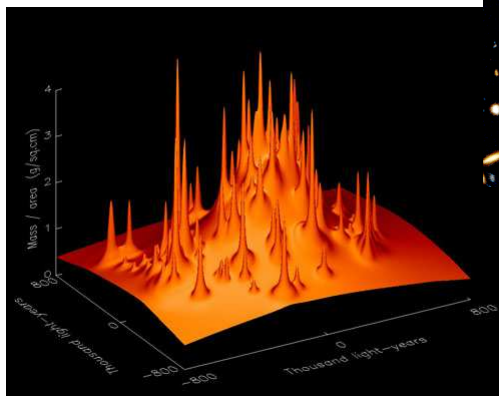


ROSAT X-ray image of Coma cluster overlaid on optical.  
MPI (ROSAT image);  
NASA/ESA/DSS2 (visible image)

5

## THE ASTROPHYSICAL EVIDENCE

### ○ Dynamics of rich clusters

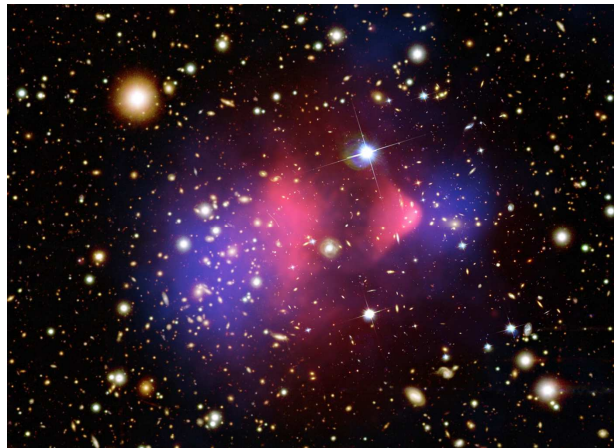


Mass map of CL0024+1654 as determined from the observed gravitational lensing.  
Tyson, Kochanski and Dell'Antonio,  
*ApJ* **498** (1998) L107

6

## THE ASTROPHYSICAL EVIDENCE: THE BULLET CLUSTER

- Mass from lens mapping (blue) follows stars not gas (red)
  - dark matter is **collisionless**

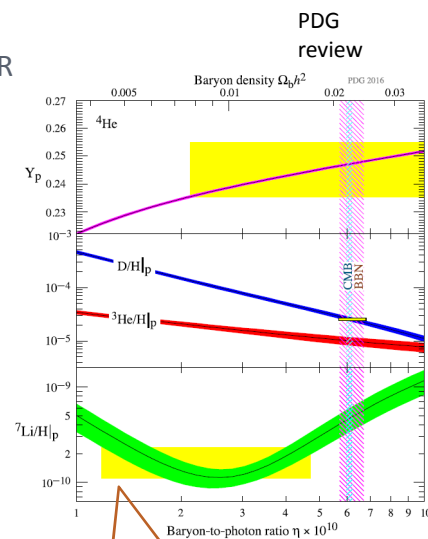


Composite Credit:  
*X-ray*: NASA/CXC/CfA/  
 M. Markevitch et al.;  
*Lensing Map*:  
 NASA/STScI; ESO WFI;  
 Magellan/U.Arizona/  
 D.Clowe et al  
*Optical*: NASA/STScI;  
 Magellan/U.Arizona/  
 D.Clowe et al.

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## NON-BARYONIC DARK MATTER

- Density of baryonic matter strongly constrained by early-universe nucleosynthesis (BBN)
  - density parameter of order 0.3 as required by data from, e.g., galaxy clusters is completely inconsistent with best fit

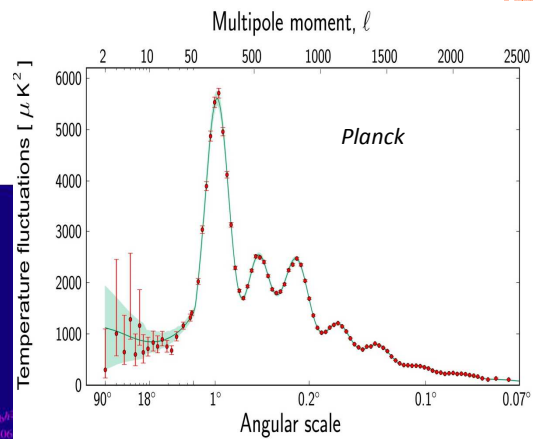
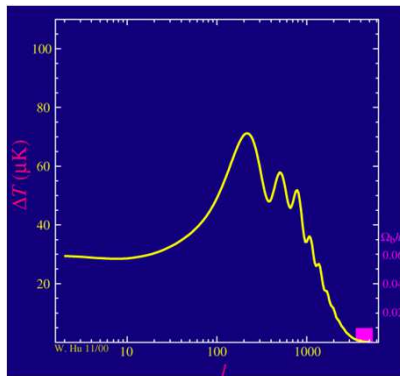


"The lithium problem":  
possible new physics??

8

## NON-BARYONIC DARK MATTER: COSMOLOGY

Ratio of odd/even peaks  
depends on  $\Omega_b$



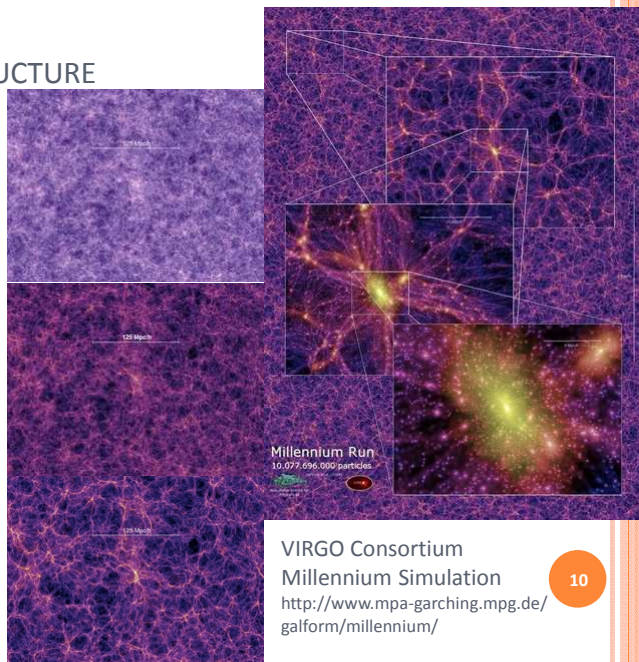
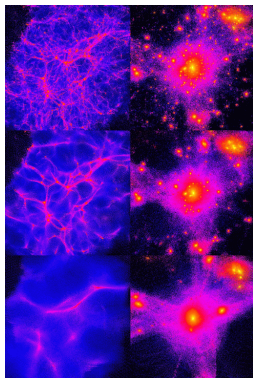
Wayne Hu

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## LARGE SCALE STRUCTURE

Relativistic (**hot**) dark  
matter makes structure  
top-down—non-relativistic  
(**cold**) bottom-up.

Real world looks like **cold**  
dark matter.

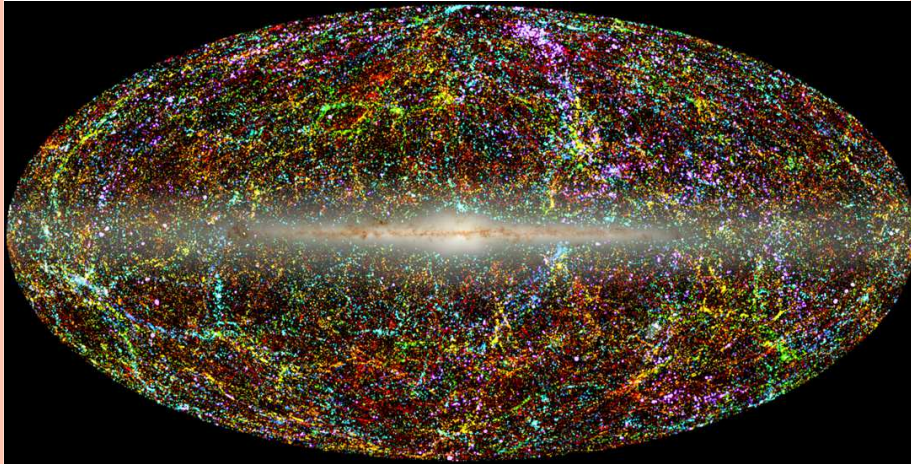


VIRGO Consortium  
Millennium Simulation  
[http://www.mpa-garching.mpg.de/  
galform/millennium/](http://www.mpa-garching.mpg.de/galform/millennium/)

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## 2MASS GALAXY SURVEY



Local galaxies ( $z < 0.1$ ; distance coded by colour, from blue to red)

Statistical studies, e.g. correlation functions, confirm visual impression that this looks much more like cold than hot dark matter

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## BRIEF SUMMARY OF ASTROPHYSICAL EVIDENCE

- Many observables concur that  $\Omega_{m0} \approx 0.3$

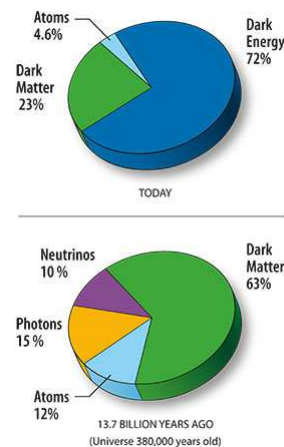
- Most of this must be non-baryonic

- BBN and CMB concur that baryonic matter contributes  $\Omega_{b0} \approx 0.05$
- Bullet Cluster mass distribution indicates that dark matter is collisionless

- No Standard Model candidate

- neutrinos are too light, and are “hot” (relativistic at decoupling)
  - hot dark matter does not reproduce observed large-scale structure

→ *BSM physics*



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## DARK MATTER

13

Astrophysical Evidence

Candidates

Detection

## DARK MATTER CANDIDATES

	WIMPs	SuperWIMPs	Light $\tilde{G}$	Hidden DM	Sterile $\nu$	Axions
Motivation	GHP	GHP	GHP/NPFP	GHP/NPFP	$\nu$ Mass	Strong CP
Naturally Correct $\Omega$	Yes	Yes	No	Possible	No	No
Production Mechanism	Freeze Out	Decay	Thermal	Various	Various	Various
Mass Range	GeV-TeV	GeV-TeV	eV-keV	GeV-TeV	keV	$\mu\text{eV}$ -meV
Temperature	Cold	Cold/Warm	Cold/Warm	Cold/Warm	Warm	Cold
Collisional				✓		
Early Universe		✓✓		✓		
Direct Detection	✓✓			✓		✓✓
Indirect Detection	✓✓	✓		✓	✓✓	
Particle Colliders	✓✓	✓✓	✓✓	✓		

GHP = Gauge Hierarchy Problem; NPFP = New Physics Flavour Problem

✓ = possible signal; ✓✓ = expected signal

Jonathan Feng, *ARAA* **48** (2010) 495 (highly recommended)

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## PARTICLE PHYSICS MOTIVATIONS

### ○ Gauge Hierarchy Problem

- in SM, loop corrections to Higgs mass give

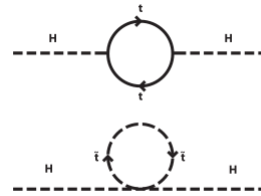
$$\Delta m_h^2 \approx \frac{\lambda^2}{16\pi^2} \int \frac{d^4 p}{p^2} \approx \frac{\lambda^2}{16\pi^2} \Lambda^2$$

and there is no obvious reason why  $\Lambda \neq M_{Pl}$

- supersymmetry fixes this by introducing a new set of loop corrections that cancel those from the SM
- new physics at TeV scale will also fix it (can set  $\Lambda \sim 1$  TeV)

### ○ New Physics Flavour Problem

- we observe conservation or near-conservation of B, L, CP
  - and do not observe flavour-changing neutral currents
- new physics has a nasty tendency to violate these
  - can require fine-tuning or new discrete symmetries, e.g. *R*-parity

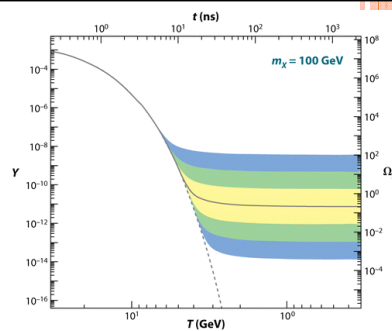



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## WIMPs

### ○ Weakly Interacting Massive Particles

- produced thermally in early universe
- annihilate as universe cools, but “freeze out” when density drops so low that annihilation no longer occurs with meaningful rate
  - “target volume” per particle in time  $\Delta t$  is  $\sigma_A v \Delta t$ , where  $\sigma_A$  is cross-section
  - so annihilation rate is  $n_f \langle \sigma_A v \rangle$  where  $n_f$  is number density
- freeze-out occurs when  $H \approx n_f \langle \sigma_A v \rangle$ , and in radiation era we have  $H \propto T^2/M_{Pl}$  (because  $\rho \propto T^4$  and  $G \propto 1/M_{Pl}^2$ )
- can estimate relic density by considering freeze-out



 Feng JL. 2010. Annu. Rev. Astron. Astrophys. 48:495–545

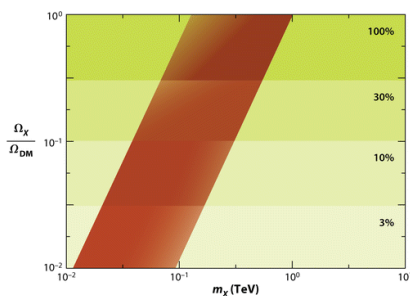
$$n_f \approx (m_\chi T_f)^{3/2} e^{-m_\chi/T_f} \approx \frac{T_f^2}{M_{Pl} \langle \sigma_A v \rangle}$$

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## WIMP RELIC DENSITY

- Converting to  $\Omega$  gives  $\Omega_X = \frac{m_X n_0}{\rho_c} \approx \frac{m_X T_0^3}{\rho_c} \frac{n_f}{T_f^3} \approx \frac{x_f T_0^3}{\rho_c M_{Pl}} \langle \sigma_A v \rangle^{-1}$   
 where  $x_f = m_X/T_f$ 
  - and typically  $\langle \sigma_A v \rangle \propto 1/m_X^2$  or  $v^2/m_X^2$  (S or P wave respectively)
- Consequence: weakly interacting massive particles with electroweak-scale masses “naturally” have reasonable relic densities



- and therefore make excellent dark matter candidates

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Feng JL. 2010.  
 Annu. Rev. Astron. Astrophys. 48:495–545

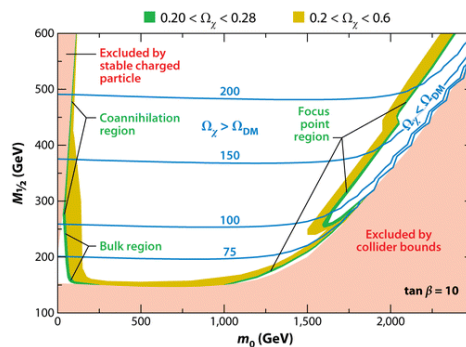
## SUPERSYMMETRIC WIMPS

- Supersymmetry solves the GHP by introducing cancelling corrections
  - predicts a complete set of new particles
  - NPFP often solved by introducing  $R$ -parity—new discrete quantum number
    - then lightest supersymmetric particle is stable
    - best DM candidate is lightest neutralino (mixed spartner of  $W^0$ ,  $B$ ,  $H$ ,  $h$ )
  - far too many free parameters in most general supersymmetric models
    - so usually consider constrained models with simplifying assumptions
    - most common constrained model: mSUGRA
      - parameters  $m_0$ ,  $M_{1/2}$ ,  $A_0$ ,  $\tan \beta$ ,  $\text{sign}(\mu)$
    - mSUGRA neutralino is probably the best studied DM candidate

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## SUSY WIMPs

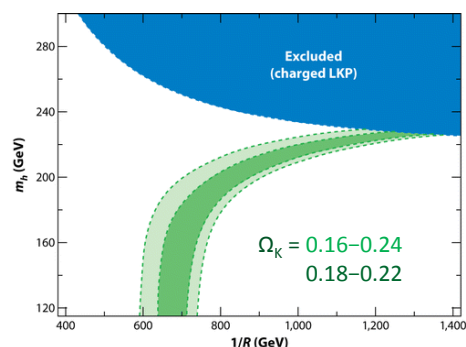
- Neutralinos are Majorana fermions and therefore self-annihilate
  - Pauli exclusion principle implies that  $\chi_1\chi_1$  annihilation prefers to go to spin 0 final state
  - $f\bar{f}$  prefers spin 1
    - therefore annihilation cross-section is suppressed
    - hence  $\Omega_\chi$  tends to be too high
    - parameter space **very** constrained by WMAP



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## KALUZA-KLEIN WIMPs

- In extra-dimension models, SM particles have partners with the *same* spin
  - “tower” of masses separated by  $R^{-1}$ , where  $R$  is size of compactified extra dimension
  - new discrete quantum number,  $K$ -parity, implies lightest KK particle is stable
    - this is the potential WIMP candidate
    - usually  $B^1$
  - annihilation not spin-suppressed (it's a boson), so preferred mass higher

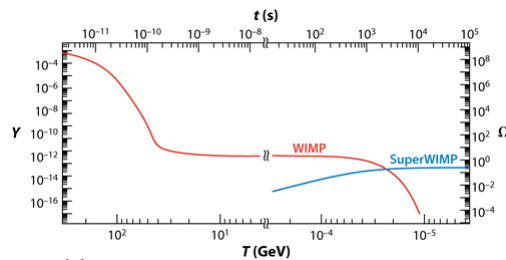


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## SUPERWIMPS

### Massive particles with superweak interactions

- produced by decay of metastable WIMP
  - because this decay is superweak, lifetime is very long ( $10^3$ – $10^7$  s)
  - WIMP may be neutralino, but could be charged particle
    - dramatic signature at LHC (stable supermassive particle)
- candidates:
  - weak-scale gravitino
  - axino
  - equivalent states in KK theories
- these particles cannot be directly detected, but indirect-detection searches and colliders may see them
  - they may also have detectable astrophysical signatures



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## LIGHT GRAVITINOS

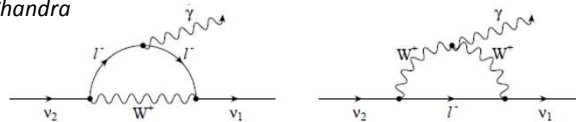
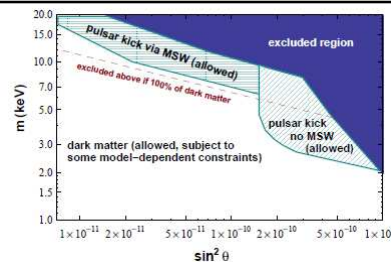
- Expected in gauge-mediated supersymmetry breaking
  - in these models gravitino has  $m < 1$  GeV
    - neutralinos decay through  $\gamma\tilde{G}$ , so cannot be dark matter
  - gravitinos themselves are possible DM candidates
    - but tend to be too light, i.e. too warm, or too abundant
    - relic density in minimal scenario is  $\Omega_{\tilde{G}} \approx 0.25 m_{\tilde{G}}/(100 \text{ eV})$ 
      - so require  $m_{\tilde{G}} < 100 \text{ eV}$  for appropriate relic density
      - but require  $m_{\tilde{G}} > 2 \text{ keV}$  for appropriate large-scale structure
  - models which avoid these problems look contrived

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Kusenko, DM10

## STERILE NEUTRINOS

- Seesaw mechanism for generating small  $\nu_L$  masses implies existence of massive right-handed sterile states
  - usually assumed that  $M_R \approx M_{\text{GUT}}$  in which case sterile neutrinos are not viable dark matter candidates
  - but smaller Yukawa couplings can combine with smaller  $M_R$  to produce observed  $\nu_L$  properties together with sterile neutrino at keV mass scale—viable dark matter candidate
    - such a sterile neutrino could also explain observed high velocities of pulsars (asymmetry in supernova explosion generating “kick”)
    - these neutrinos are not entirely stable:  $\tau \gg 1/H_0$ , but they do decay and can generate X-rays via loop diagrams—therefore potentially detectable by, e.g., *Chandra*



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## STERILE NEUTRINOS

- Production mechanisms
  - oscillation at  $T \approx 100$  MeV
    - $\Omega_\nu \propto \sin^2 2\theta m^{1.8}$  from numerical studies
    - always present: requires small mass and very small mixing angle
      - not theoretically motivated: some fine tuning therefore required
  - resonant neutrino oscillations
    - if universe has significant lepton number asymmetry,  $L > 0$
  - decays of heavy particles
    - e.g. singlet Higgs driving sterile neutrino mass term
- Observational constraints
  - X-ray background
  - presence of small-scale structure
    - sterile neutrinos are “warm dark matter” with Mpc free-streaming

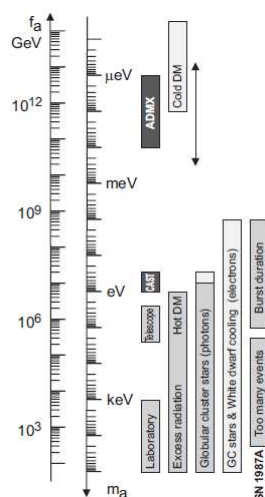
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## AXIONS

- Introduced to solve the “strong CP problem”
  - SM Lagrangian includes CP-violating term which should contribute to, e.g., neutron electric dipole moment
    - neutron doesn’t appear to *have* an EDM ( $<3 \times 10^{-26}$  e cm, cf. naïve expectation of  $10^{-16}$ ) so this term is strongly suppressed
  - introduce new pseudoscalar field to kill this term (Peccei-Quinn mechanism)
    - result is an associated pseudoscalar boson, the axion
- Axions are extremely light ( $<10$  meV), but are cold dark matter
  - not produced thermally, but via phase transition in very early universe
    - if this occurs before inflation, visible universe is all in single domain
    - if after inflation, there are many domains, and topological defects such as axion domain walls and axionic strings may occur

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## AXIONS

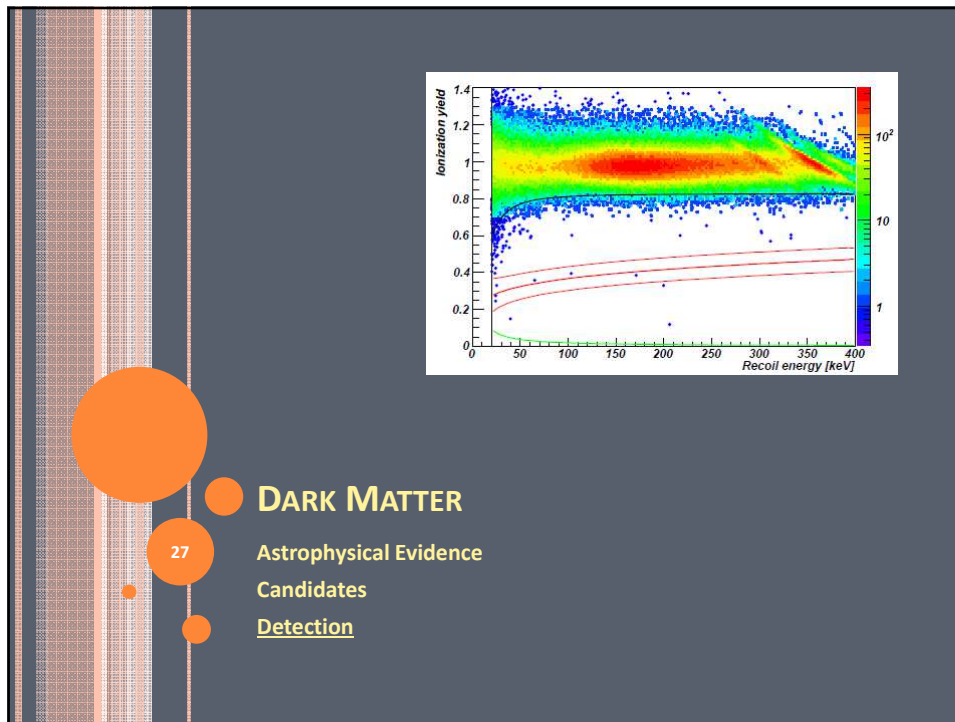


Georg Raffelt, hep-ph/0611350v1

- Axion mass is  $m_a \approx 6 \mu\text{eV} \times f_a / (10^{12} \text{ GeV})$  where  $f_a$  is the unknown mass scale of the PQ mechanism
- Calculated relic density is  $\Omega_a \approx 0.4 \vartheta^2 (f_a / 10^{12} \text{ GeV})^{1.18}$  where  $\vartheta$  is initial vacuum misalignment
  - so need  $f_a < 10^{12} \text{ GeV}$  to avoid overclosing universe
  - astrophysical constraints require  $f_a > 10^9 \text{ GeV}$
  - therefore  $6 \mu\text{eV} < m_a < 6 \text{ meV}$

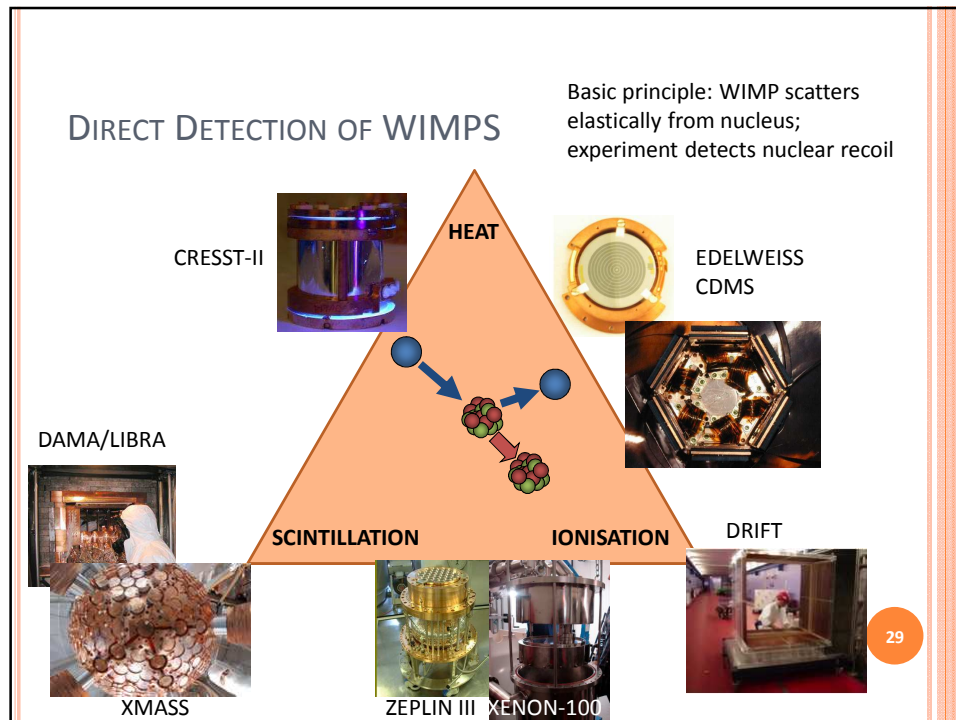
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## DETECTION OF DARK MATTER CANDIDATES

- Direct detection
  - dark matter particle interacts in your detector and you observe it
- Indirect detection
  - you detect its decay/annihilation products or other associated phenomena
- Collider phenomenology
  - it can be produced at, say, LHC and has a detectable signature
- Cosmology
  - it has a noticeable and characteristic impact on BBN or CMB
- Focus here on best studied candidates—WIMPs and axions



## DIRECT DETECTION OF WIMPS

- **Backgrounds**
  - cosmics and radioactive nuclei (especially radon)
    - use deep site and radiopure materials
    - use discriminators to separate signal and background
- **Time variation**
  - expect annual variation caused by Earth's and Sun's orbital motion
    - small effect, ~7%
    - basis of claimed signal by DAMA experiment
  - much stronger diurnal variation caused by changing orientation of Earth
    - "smoking gun", but requires directional detector
    - current directional detector, DRIFT, has rather small target mass (being gaseous)—hence not at leading edge of sensitivity

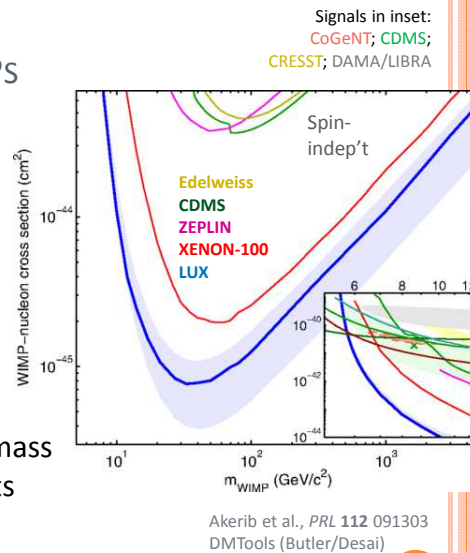
CDMS-II, *PRL*  
106 (2011)  
131302

ZEPLIN-II, *Astropart. Phys.*  
28 (2007) 287

30

## DIRECT DETECTION OF WIMPS

- Interaction with nuclei can be spin-independent or spin-dependent
  - spin-dependent interactions require nucleus with net spin
  - most direct detection experiments focus on SI, and limits are much better in this case
- Some claimed signals at low mass inconsistent with others' limits
  - requires very low mass and high cross-section
    - if real, may point to a non-supersymmetric DM candidate



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## INDIRECT DETECTION OF WIMPS

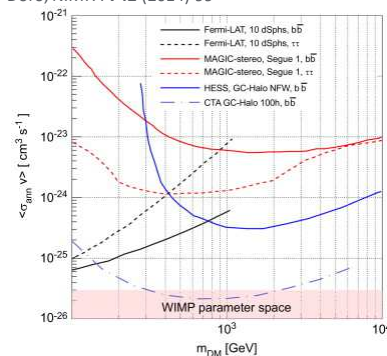
- After freeze-out, neutralino self-annihilation is negligible in universe at large
  - but neutralinos can be captured by repeated scattering in massive bodies, e.g. Sun, and this will produce a significant annihilation rate
    - number of captured neutralinos  $N = C - AN^2$  where  $C$  is capture rate and  $A$  is  $\langle \sigma_A v \rangle$  per volume
    - if steady state reached, annihilation rate is just  $C/2$ , therefore determined by scattering cross-section
  - annihilation channels include  $W^+W^-$ ,  $b\bar{b}$ ,  $\tau^+\tau^-$ , etc. which produce secondary neutrinos
    - these escape the massive object and are detectable by neutrino telescopes

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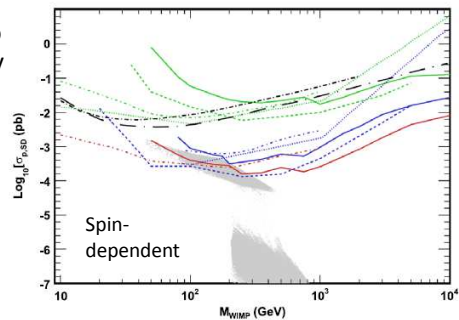
## INDIRECT DETECTION OF WIMPs

- Relatively high threshold of neutrino telescopes implies greater sensitivity to “hard” neutrinos, e.g. from WW
- Also possible that neutralinos might collect near Galactic centre

Doro, *NIMPA 742* (2014) 99



— ANTARES; ... Super-K; --- IceCube  
( $b\bar{b}$ ;  $W\bar{W}$ ;  $\tau\bar{\tau}$ )  
--- SIMPLE; --- COUPP



Zornoza & Lambard, *NIMPA 742* (2014) 173

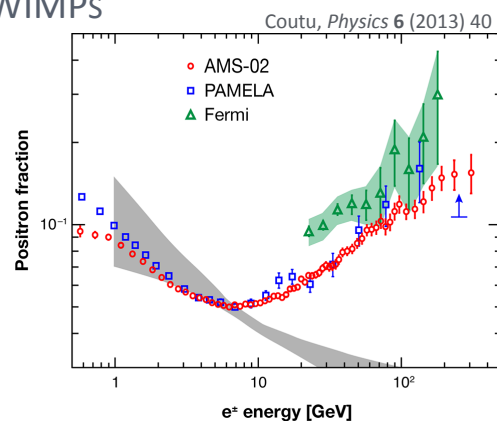
- in this region other annihilation products, e.g.  $\gamma$ -rays, could escape
- Limits from Fermi-LAT, MAGIC, H.E.S.S. not yet reaching expected signal, but CTA could come close

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## INDIRECT DETECTION OF WIMPs

- Various experiments see positron fraction in cosmic rays rising with energy (unexpectedly)

- This *could* be a signal of  $\chi\chi$  annihilation
- But there are also potential astrophysical sources, e.g. pulsar wind nebulae, microquasars
- And there are no signals for  $\gamma$ -ray excess, which we might expect



Coutu, *Physics 6* (2013) 40

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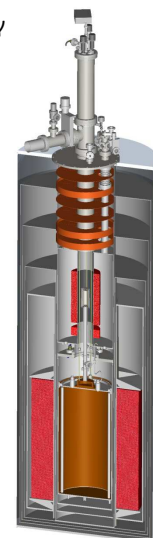
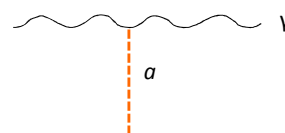
## LHC DETECTION OF WIMPs AND SWIMPS

- WIMPs show up at LHC through missing-energy signature
  - note: not immediate proof of dark-matter status
    - long-lived but not stable neutral particle would have this signature but would not be DM candidate
    - need to constrain properties enough to calculate expected relic density if particle *is* stable, then check consistency
- SuperWIMP parents could also be detected
  - if charged these would be spectacular, because of extremely long lifetime
    - very heavy particle exits detector without decaying
      - if seen, could in principle be trapped in external water tanks, or even dug out of cavern walls (Feng: “new meaning to the phrase ‘data mining’”)
  - if neutral, hard to tell from WIMP proper
    - but mismatch in relic density, or conflict with direct detection, possible clues

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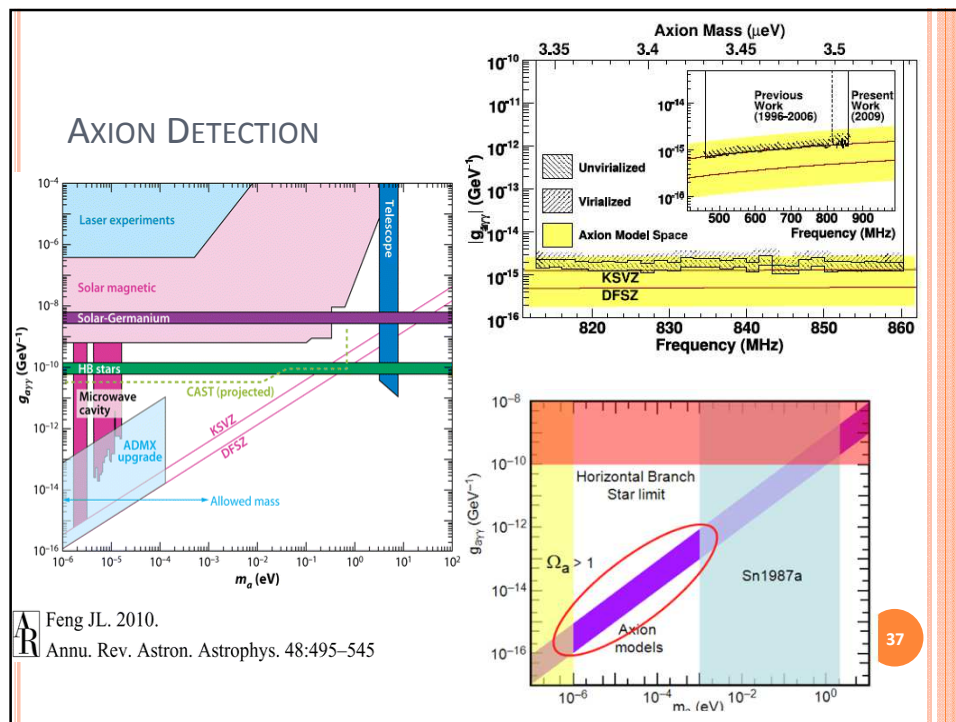
## AXION DETECTION

- Axions couple (unenthusiastically) to photons via
 
$$\mathcal{L}_{a\gamma\gamma} = -g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$
  - they can therefore be detected using Primakoff effect (resonant conversion of axion to photon in magnetic field)
  - ADMX experiment uses very high  $Q$  resonant cavity in superconducting magnet to look for excess power
- this is a scanning experiment:
  - need to adjust resonant frequency to “see” specific mass (*very tedious*)
- alternative: look for axions produced in Sun (CAST)
  - non-scanning, but less sensitive



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## DARK MATTER: SUMMARY

- Astrophysical evidence for dark matter is consistent and compelling
  - not an unfalsifiable theory—for example, severe conflict between BBN and WMAP on  $\Omega_b$  might have scuppered it
- Particle physics candidates are many and varied
  - and in many cases are not *ad hoc* inventions, but have strong independent motivation from within particle physics
- Unambiguous detection is possible for several candidates, but will need careful confirmation
  - interdisciplinary approaches combining direct detection, indirect detection, conventional high-energy physics and astrophysics may well be required

# THE END



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